

Electrons, phonons and screening

Matthieu J Verstraete, ULiege Belgium

Exciting NEWS 2021

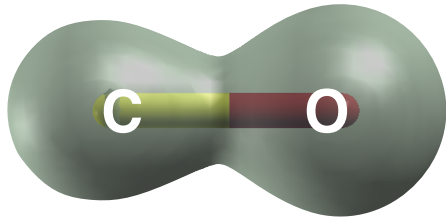
~~HU~~Berlin July 2020

~~Riga~~ June 2021

on Zoom June 23rd 2021

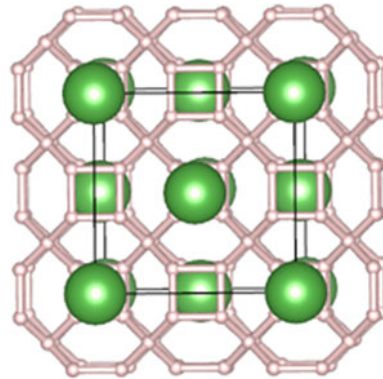
First principles quantum simulations

$$\mathcal{H}\Psi = E\Psi$$

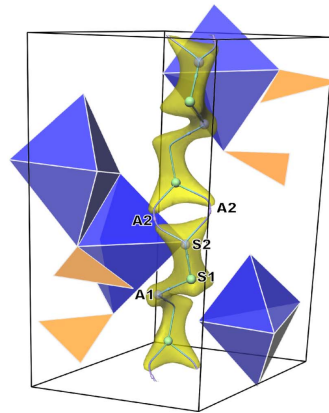


Calculate:

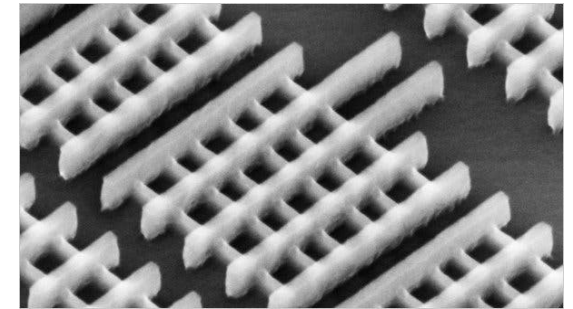
- Structure
- Bonding
- Magnetism
- **Thermodynamics**
- **Kinetics**
- **Superconductivity**
- **Transport**



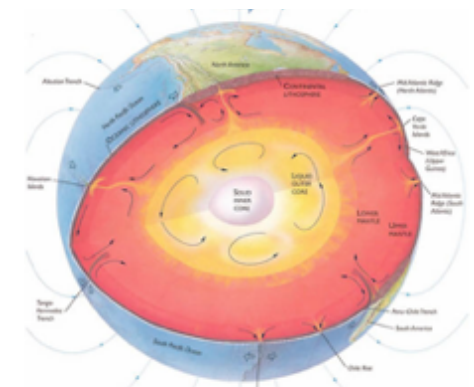
RT superconductor
Drozdov *Nature* 569 528 (2019)



Battery anodes
Kim J. *ECS* 158 A309 (2011)

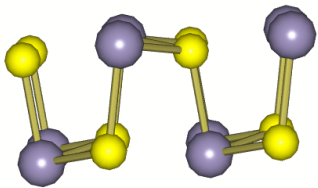
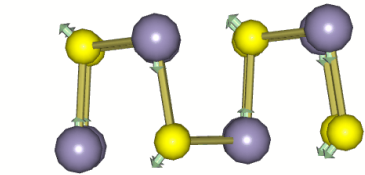


Intel HfO₂ dielectric



Earth's core
Gomi Verstraete et al.
PEPI 224 88 (2013)

DPFT refresher



Start from DFT
what moves?

- direct
- induced (SCF)
- Pulay/basis set

$$H_k^{(0)} \left| \psi_{ik}^{(0)} \right\rangle = \epsilon_{ik}^{(0)} \left| \psi_{ik}^{(0)} \right\rangle$$

$$H^{(0)} = T + V_{loc+NL}(R) + V_{Hxc}[n(R)]$$

↑
↑

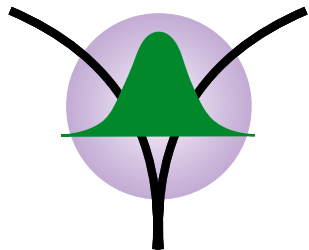
$$H_k^\alpha \left| \psi_{ik}^\alpha \right\rangle = \epsilon_{ik}^\alpha \left| \psi_{ik}^\alpha \right\rangle$$

Pertrub 1st order

+ choose gauge

$$\left(H_k^{(0)} - \epsilon_{ik}^{(0)} \right) \left| \psi_{ik}^\alpha \right\rangle = - \left(H_k^\alpha - \epsilon_k^{(\alpha)} \right) \left| \psi_{ik}^{(0)} \right\rangle$$

$$\left\langle \psi_j^{(0)} \left| \psi_i^\alpha \right\rangle = 0 \quad \forall j$$



See lecture on phonons by Fabio Caruso:
<https://www2.physik.hu-berlin.de/how-exciting/talk-fabio-caruso.mp4>

EPC refresher

Sternheimer:

- variational SCF
- project on unoccupied states

$$P_{ck} \left(H_k^{(0)} - \varepsilon_{ik}^{(0)} \right) P_{ck} \left| \psi_{ik}^\alpha \right\rangle = - P_{ck} H_k^\alpha \left| \psi_{ik}^{(0)} \right\rangle$$

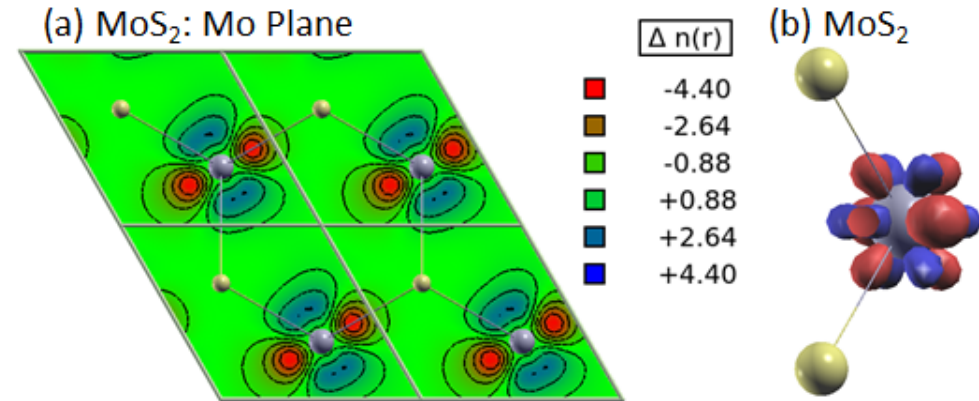
Advantages:

- "easier" than Kohn Sham
- also q dependent (any q!)

- extract coupling:
$$g_{kij}^{\alpha q} = \left\langle \psi_{ik+q}^{(0)} \left| H_k^{\alpha q} \right| \psi_{jk}^{(0)} \right\rangle$$

See the lecture on EPC by Sebastian Tillack:

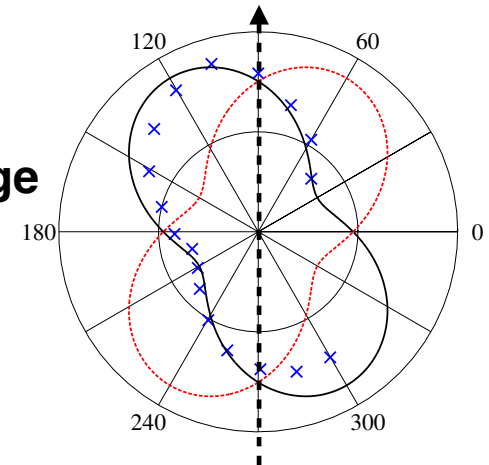
<https://www2.physik.hu-berlin.de/how-exciting/talk-sebastian-tillack.mp4>



Pike Verstraete PRB 95 201106 (2017)

Flipped sign for Born Effective Charge

Signature in Polarized Raman

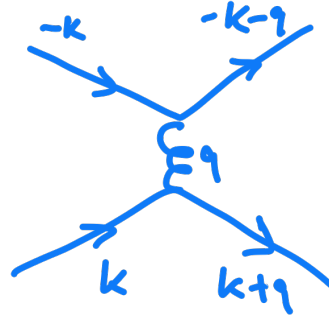


Now what does it do?

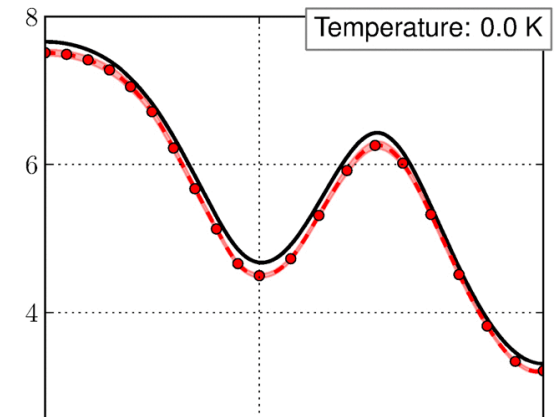
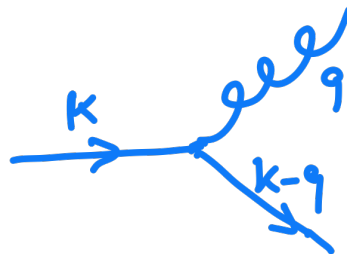
Electron energy renormalization
(AHC, polarons)



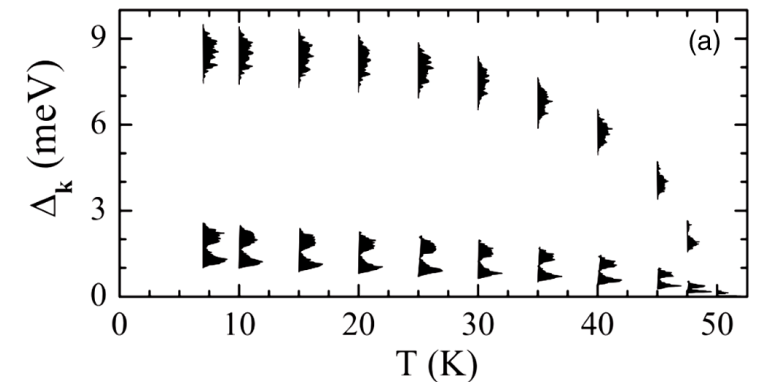
Electron pairing
(superconductivity)



Electron scattering
(resistance, Seebeck)

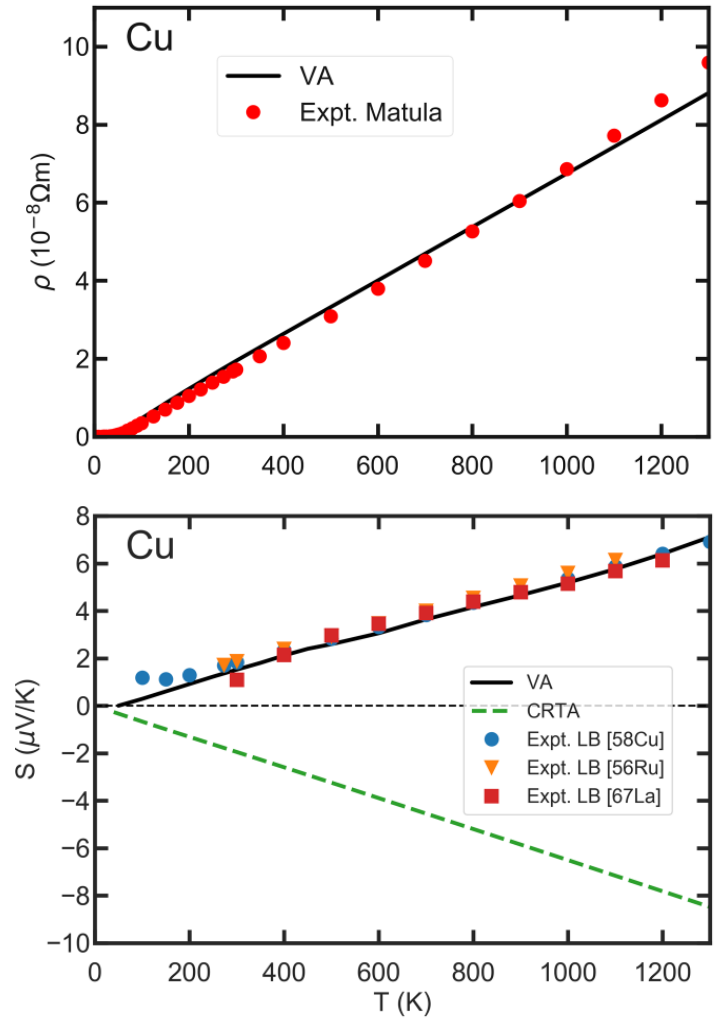
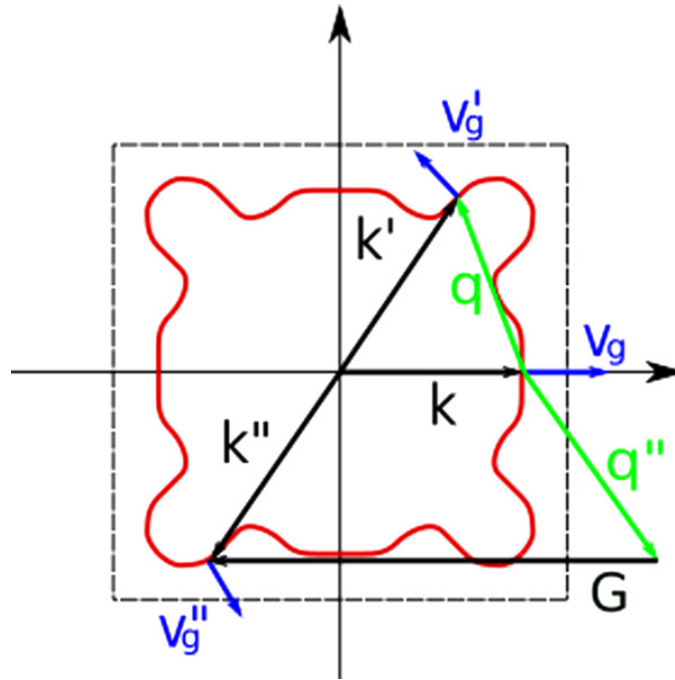
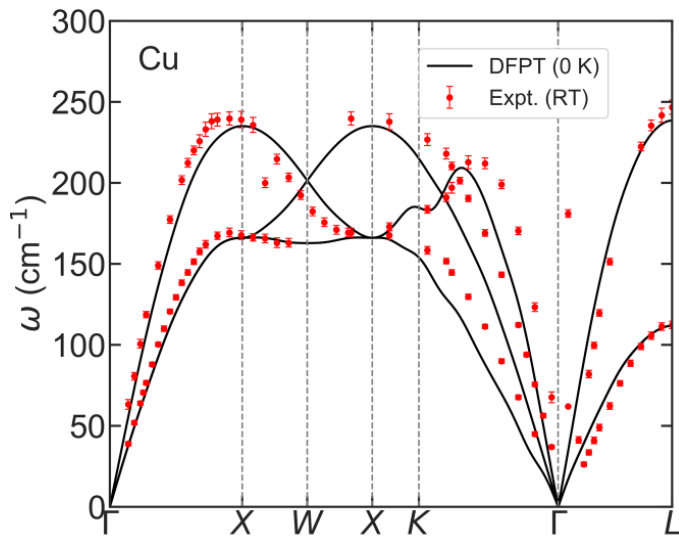
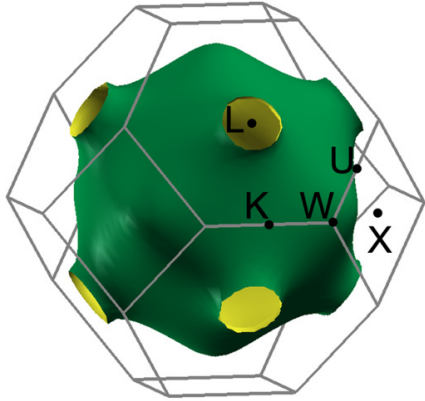


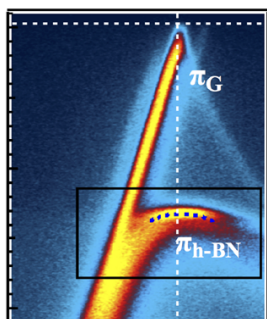
Poncé J Chem Phys 143 102813 (2015)



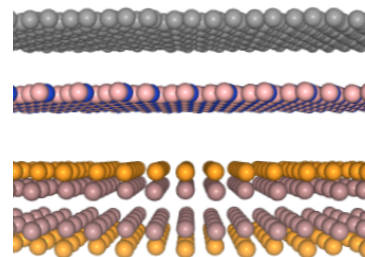
Margine Phys. Rev. B 87 024505 (2013)

Transport, thermoelectricity

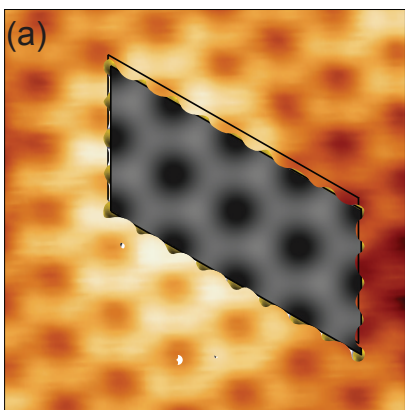




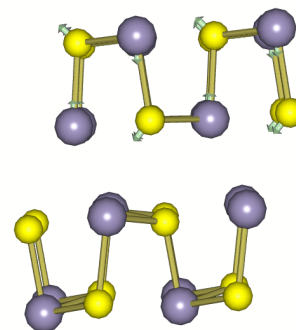
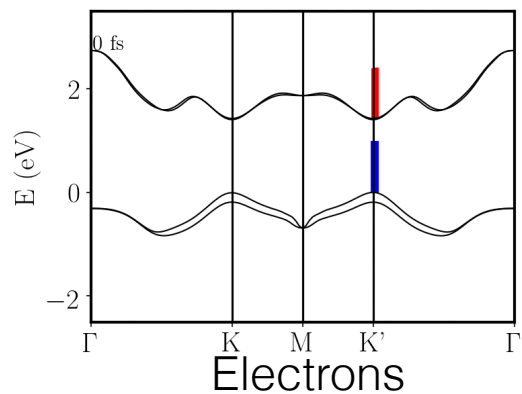
Spectroscopy



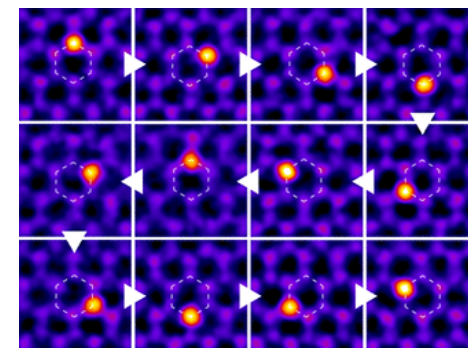
Transport



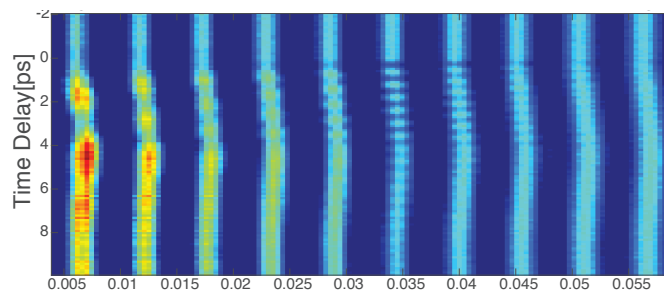
Microscopy



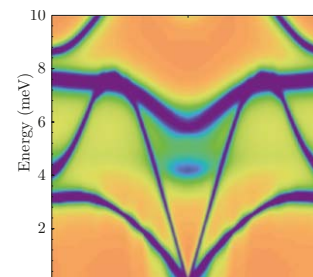
Phonons



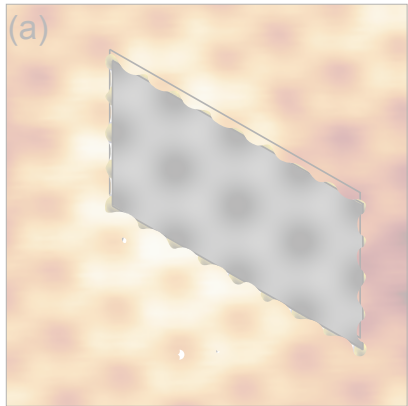
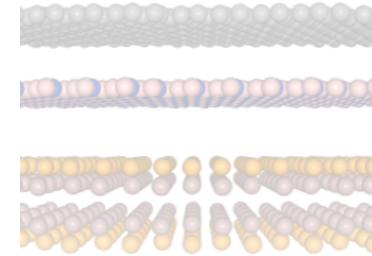
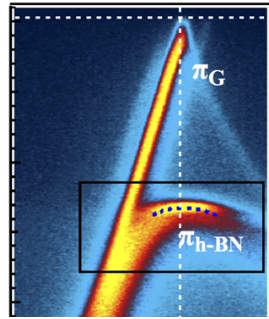
Defects



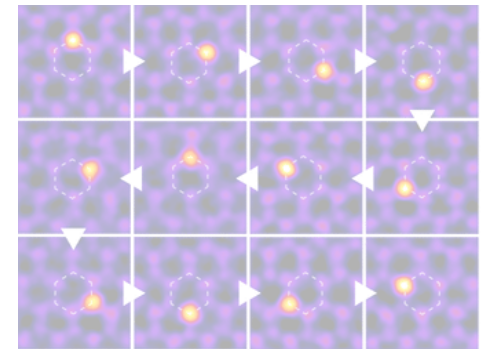
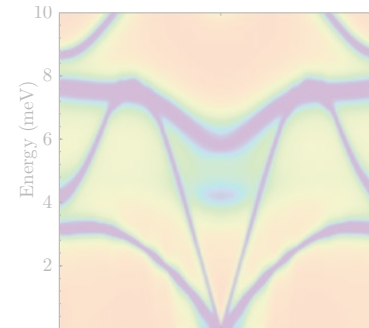
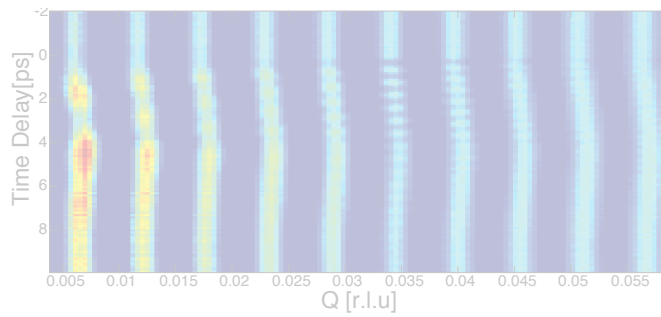
Ultrafast (FEL, pump/probe)



Synchrotron neutrons

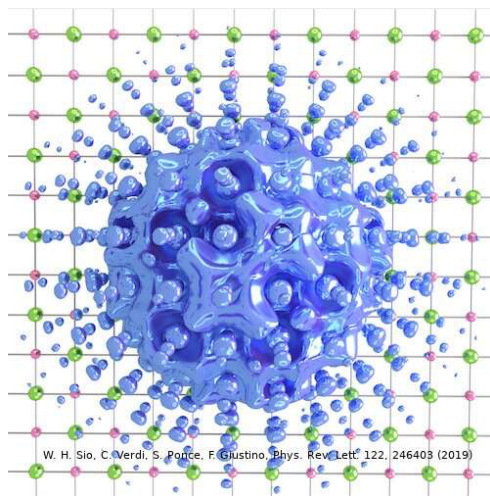


High Throughput Polarons



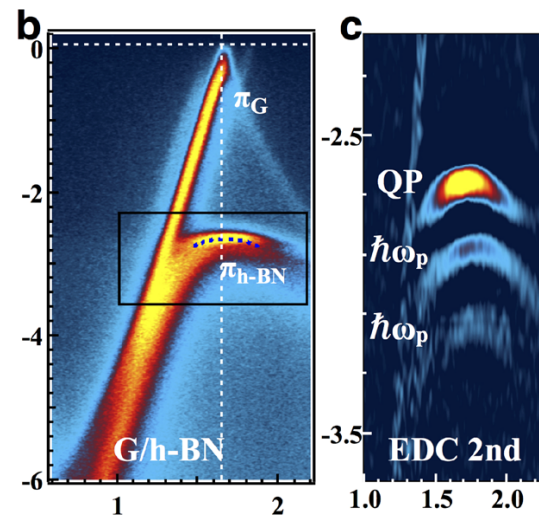
Polaron overview

What is a polaron anyway?

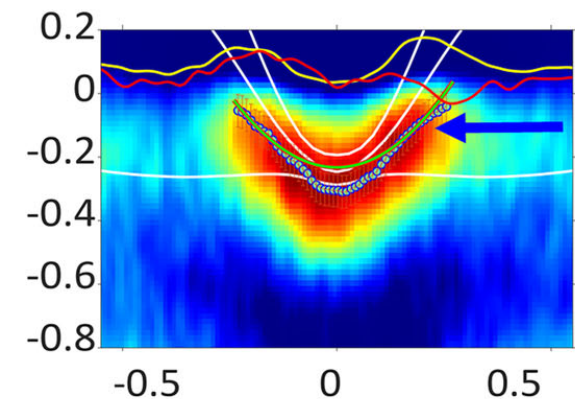


Sio PhysRevLett (2019)

Signatures in optics, transport, superconductivity...



Chen Nano Lett (2018)



CaCeMnO₃

Husanu Nat Comms (2020)

- How common are different types?
- How far can we get with simple models?

See previous lecture on polarons and ARPES by Carla Verdi!

Fröhlich model

1 electron 1 polar phonon

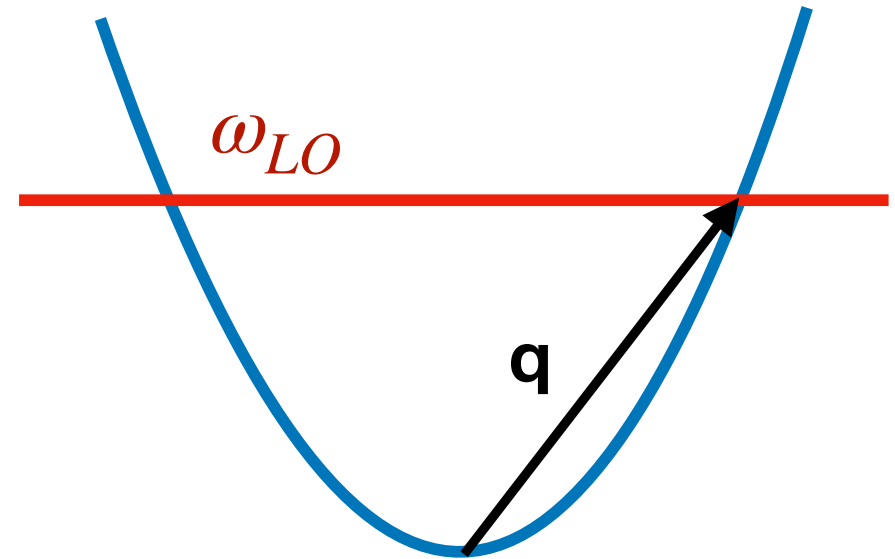
Long range electrostatic coupling

Single free parameter α

Binding energy $E_P \simeq -\alpha\omega_{LO}$

aka Zero Point Renormalization (ZPR)

Perturbation theory catastrophe at $\alpha=6$



Fröhlich Adv. Phys. (1954)

Mishchenko PRB (2000)

Story PRB (2014)

Sjakste PRB (2015)

Verdi PRL (2015)

$$\alpha = \left(\frac{1}{\epsilon^\infty} - \frac{1}{\epsilon^0} \right) \sqrt{\frac{m^*}{2\omega_{LO}}} \quad \frac{m_p^*}{m^*} \approx \left(1 - \frac{\alpha}{6} \right)^{-1}$$

Methods

Descriptors: m^* ϵ^∞ ϵ^0 ω_0

Materials Project databases:

- Intersection of > 1039 materials:
- 9000 m^* : [F. Ricci Sci Data \(2017\)](#)
- 1500 phonons: [G. Petretto Sci Data \(2018\)](#)

Criteria:

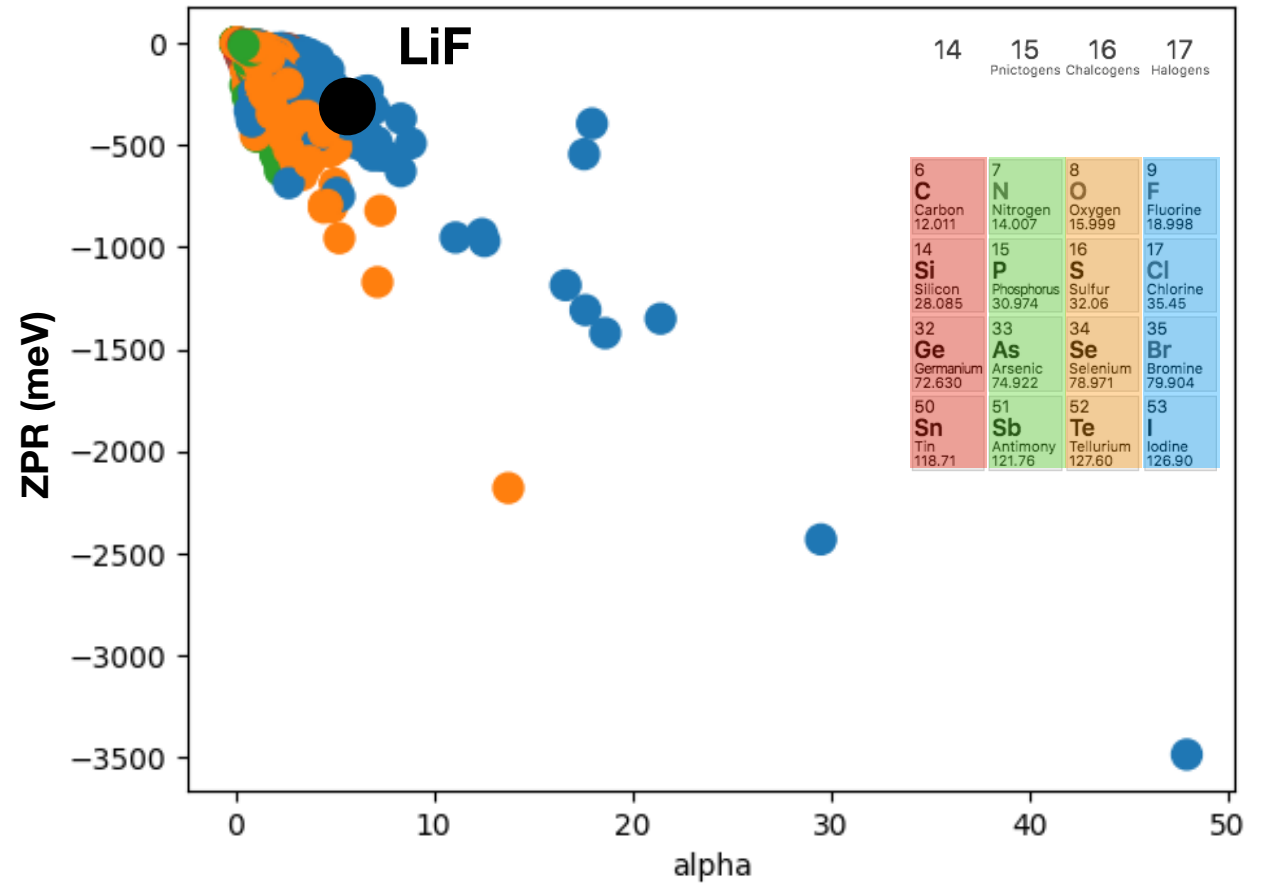
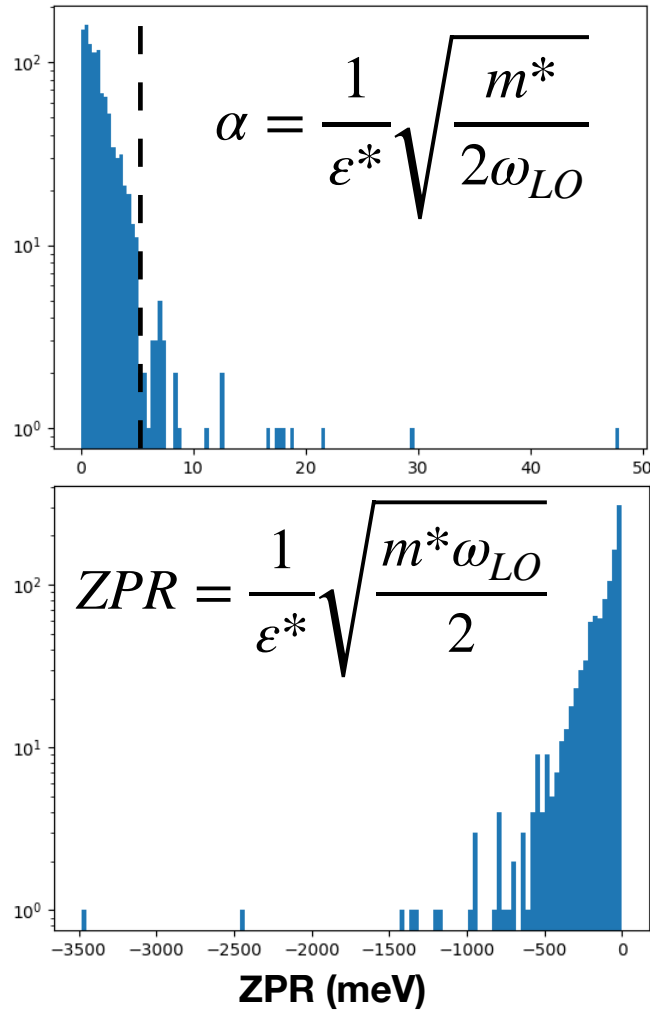
- ICSD stable 3D, nonmagnetic, insulators
- 2-5 elements per unit cell

Conduction band minimum (also valence)

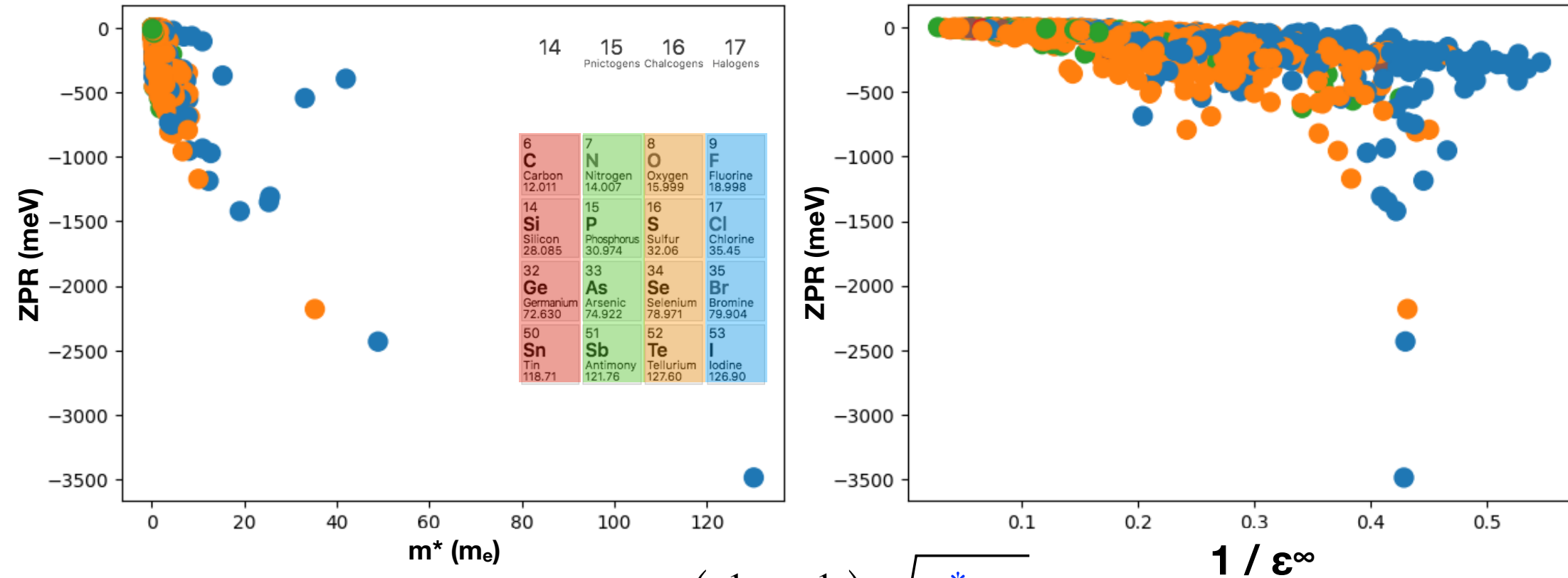


$$\alpha = \left(\frac{1}{\epsilon^\infty} - \frac{1}{\epsilon^0} \right) \sqrt{\frac{m^*}{2\omega_{LO}}}$$

Distribution of alpha and ZPR

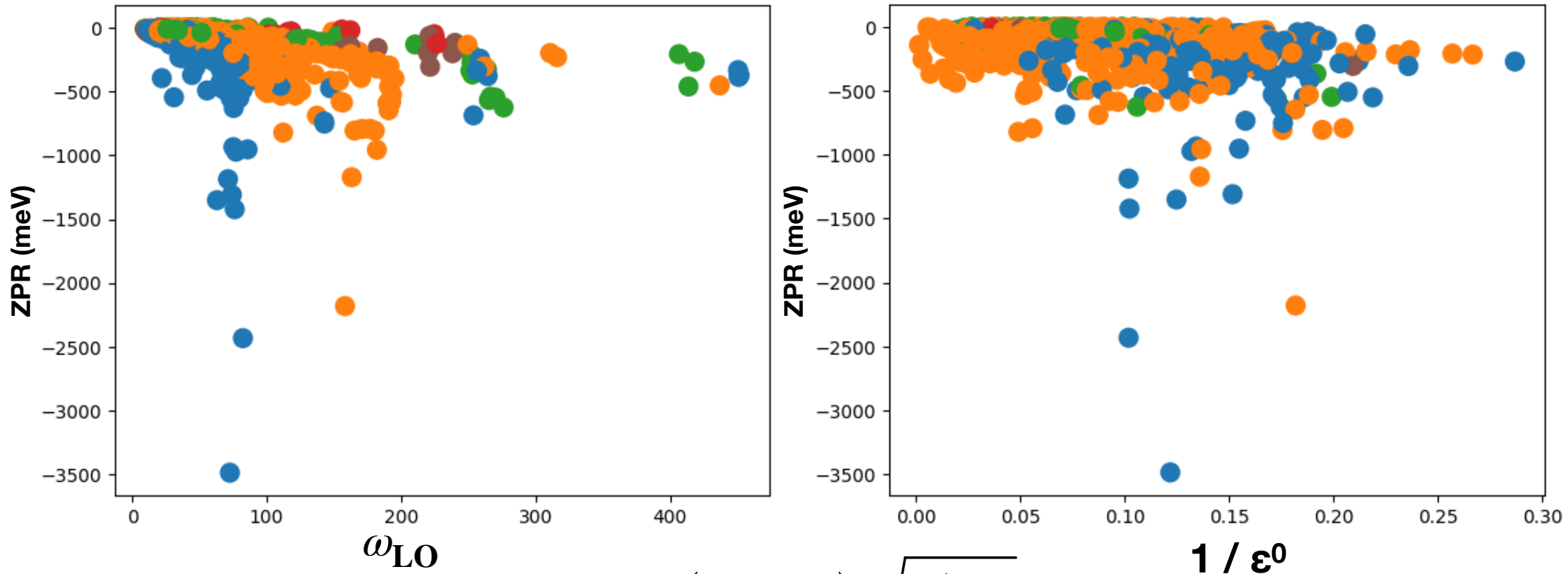


Electron effects



$$ZPR = \left(\frac{1}{\epsilon^\infty} - \frac{1}{\epsilon^0} \right) \sqrt{\frac{m^* \omega_{LO}}{2}}$$

Phonon effects



$$ZPR = \left(\frac{1}{\epsilon^\infty} - \frac{1}{\epsilon^0} \right) \sqrt{\frac{m^* \omega_{LO}}{2}}$$

Beyond Fröhlich - Generalization

Degenerate bands

All phonon modes

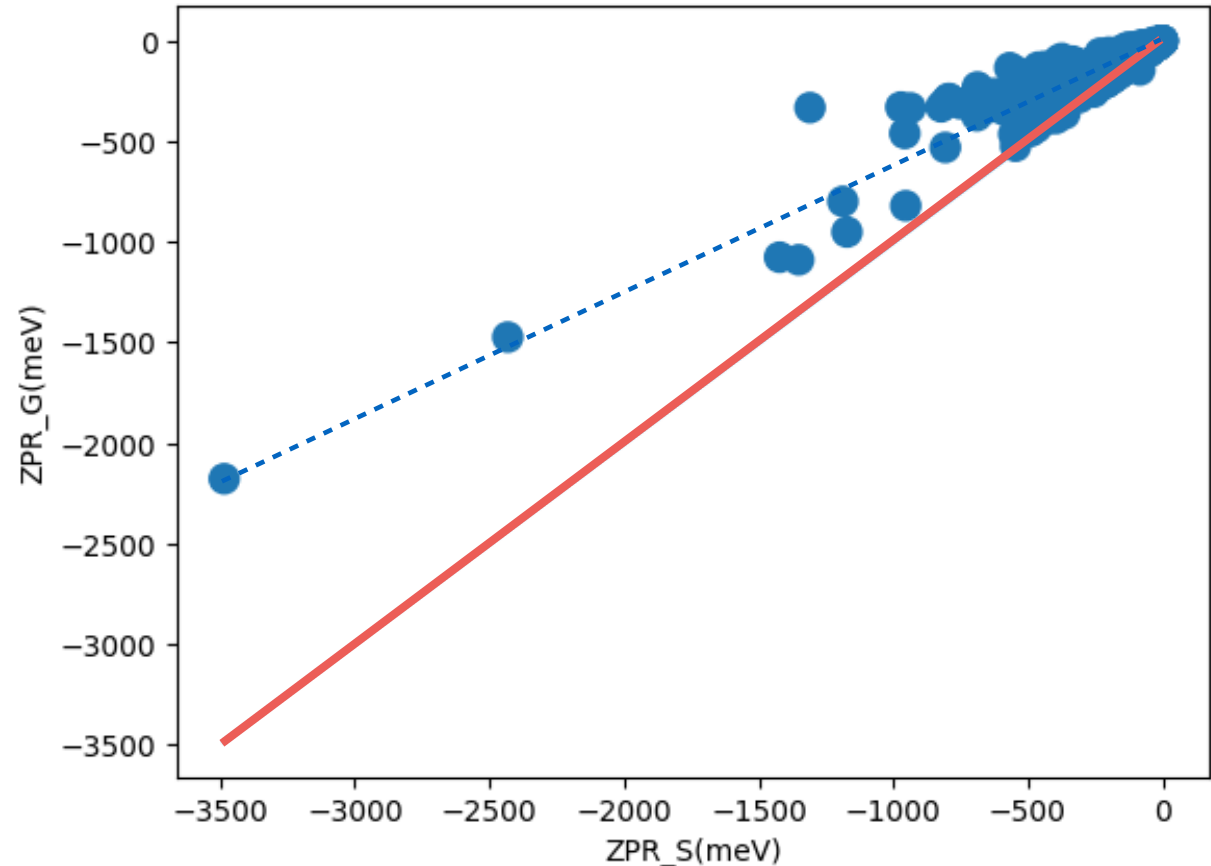
Directional m^*

Directional variation of ϵ and ω

$$\alpha_j = \int d\hat{q} \frac{\langle (m_n^*(\hat{q}))^{1/2} \rangle_n}{\epsilon_j^*(\hat{q}) \sqrt{2\omega_j(\hat{q})}}$$

A. Miglio npj Comp Mater (2020)

Why does Frohlich work so well?



Beyond Fröhlich - validation



How far off are the Fröhlich models?

Which modes contribute?

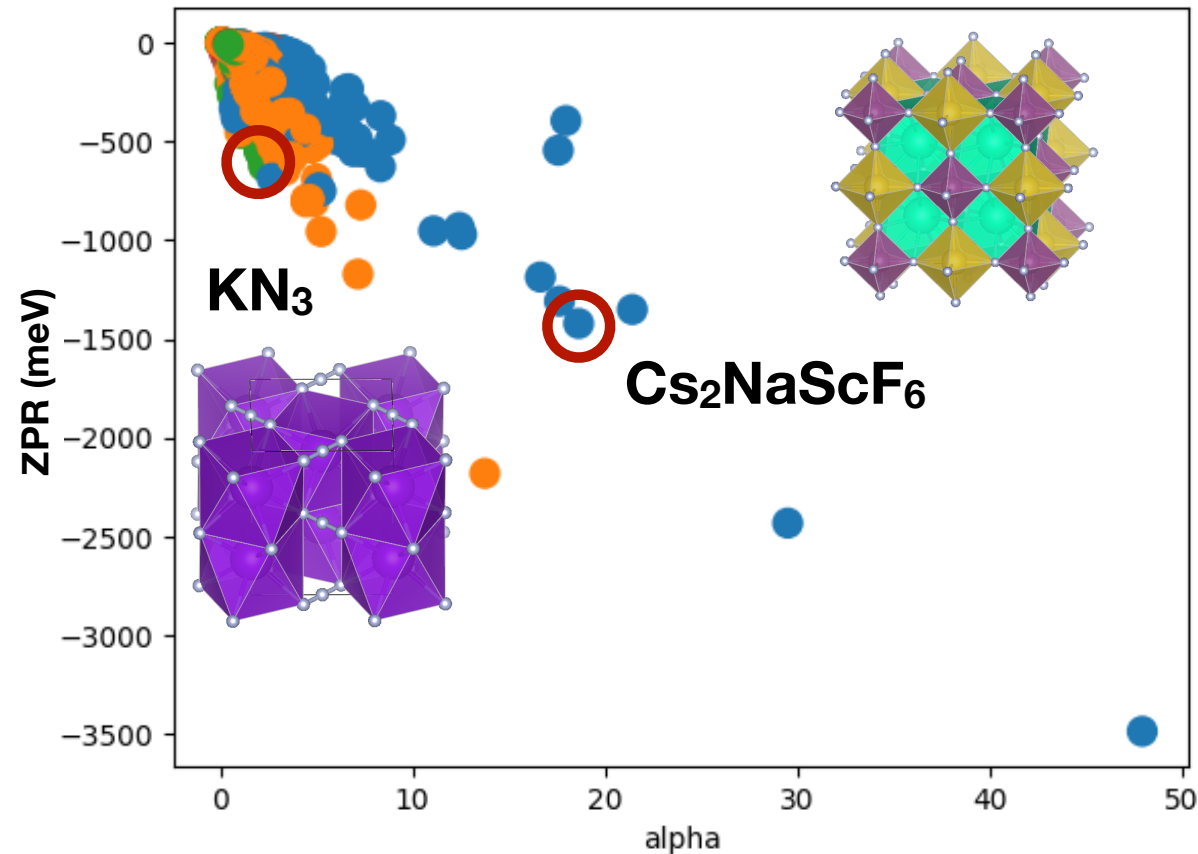
Calculate ZPR with AHC theory

DFPT Phonons, m^*

Nery PRB (2018)

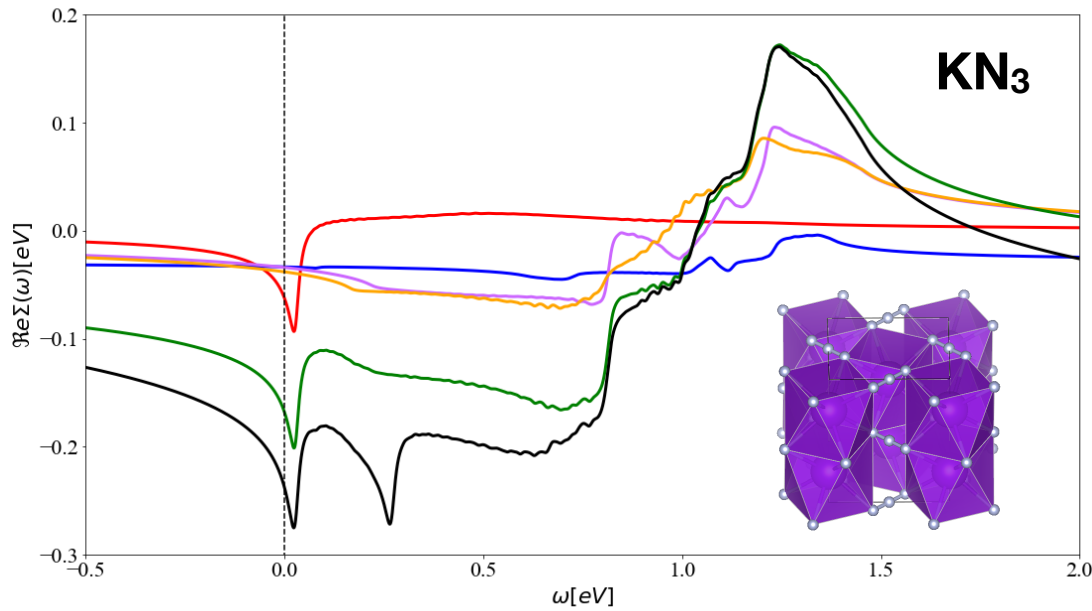
KN_3 : large ZPR small alpha

$\text{Cs}_2\text{NaScF}_6$: large ZPR large alpha

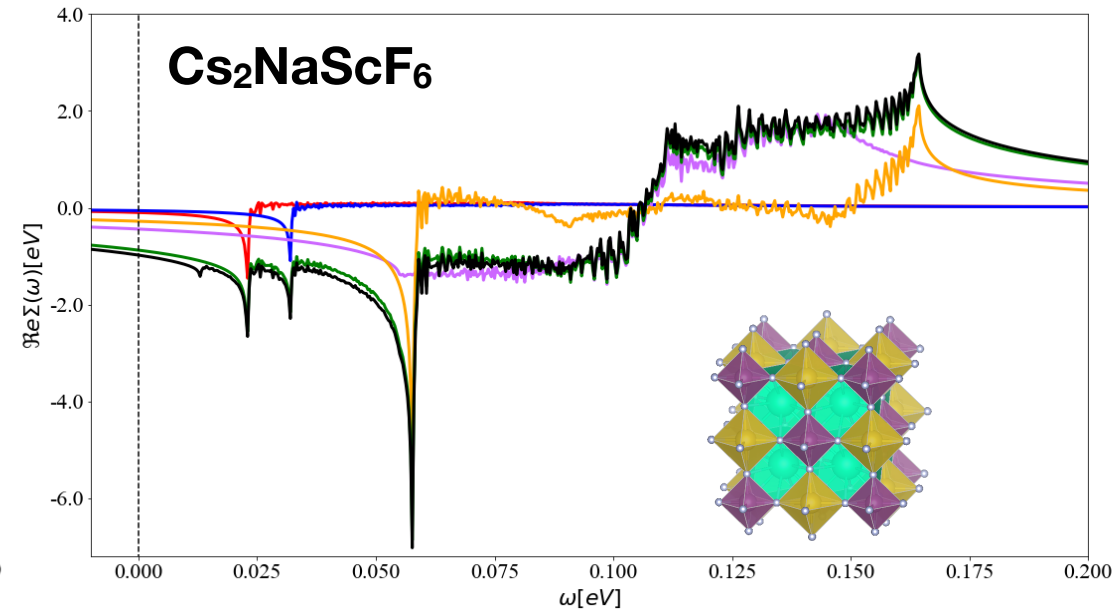


Beyond Fröhlich - validation

$$\Sigma_{\mathbf{k}nj}^{\text{Fan}}(\omega) = \frac{1}{N_q} \sum_{\mathbf{q}}^{\text{BZ}} \sum_{n'} \left| \left\langle \mathbf{k} + \mathbf{q}n' \left| H_{\mathbf{q}j}^{(1)} \right| \mathbf{k}n \right\rangle \right|^2 \times \left[\frac{1 - f_{\mathbf{k}+\mathbf{q}n'}}{\omega - \epsilon_{\mathbf{k}+\mathbf{q}n'} - \omega_{\mathbf{q}j} + i\eta} + \frac{f_{\mathbf{k}+\mathbf{q}n'}}{\omega - \epsilon_{\mathbf{k}+\mathbf{q}n'} + \omega_{\mathbf{q}j} + i\eta} \right]$$



sFr(DB) sFr(Abinit) gFr(DB) **AHC**
 -542 -330 -153 **-245**

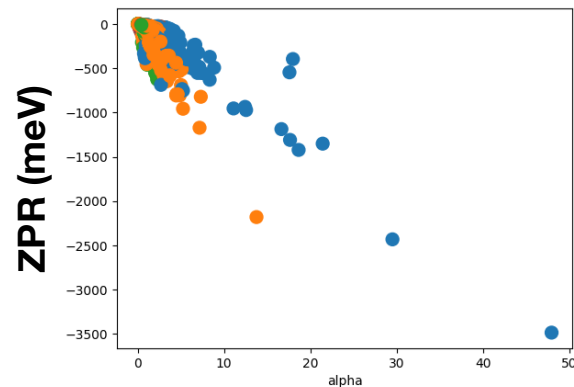


sFr(DB) sFr(Abinit) gFr(DB) **AHC**
 -1351 -1241 -1089 **-966**

Polaron takeaway

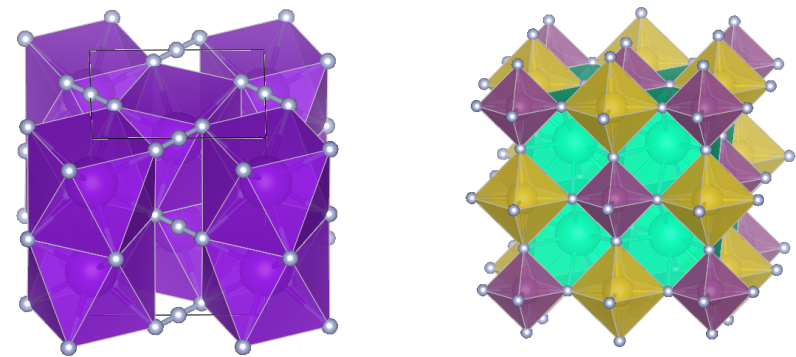
High throughput Fröhlich

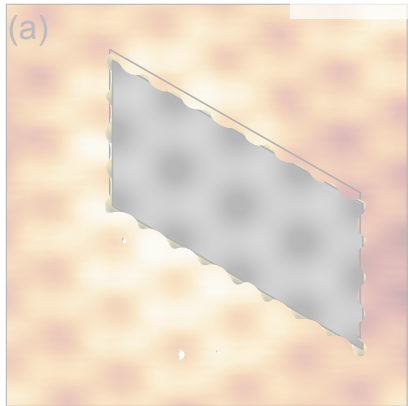
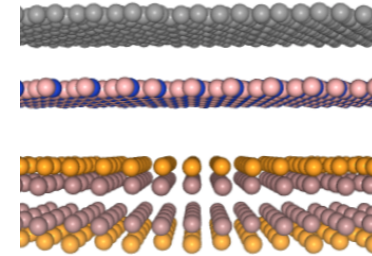
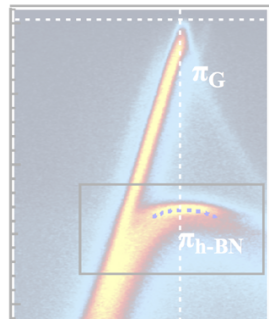
- Clustering by element period/ionicity
- Monotonic electron contribution
- Competing constraints in ϵ_0 ω_{LO} ϵ_∞ m^*
- Valid predictor for order of magnitude



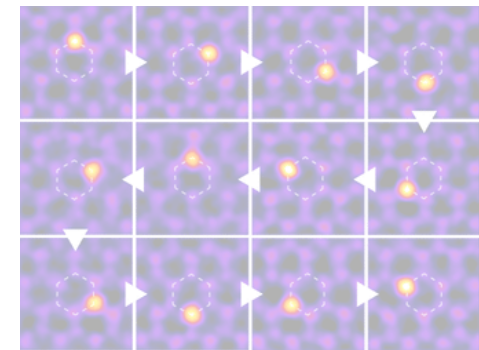
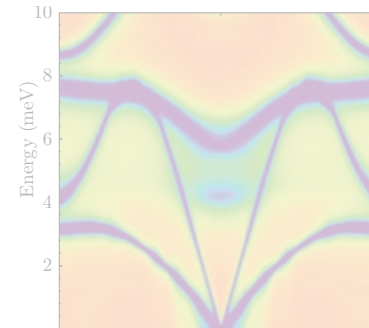
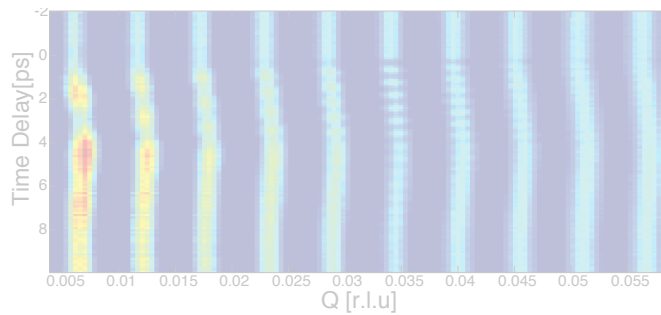
Comparison with full DFPT + AHC theory:

- Many phonons intervene
- Highest LO does not dominate
- Fröhlich ZPR 30-50% overestimation





Dielectric nano-engineering



2D Electronics

Goal: maximize electron mobility

Also for opto- or spintronics

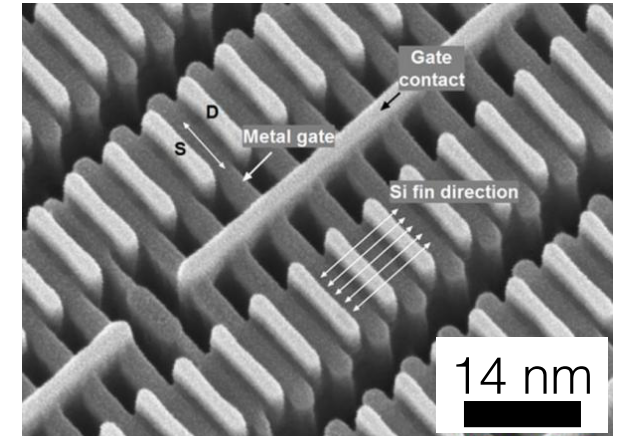
Issues for 3D materials @ nanoscale:

leakage, heat, fabrication, cost

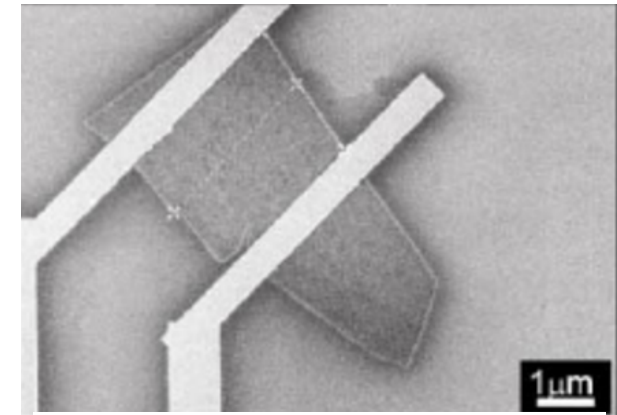
What changes in 2D?

- + low mass
- + low operational voltage
- + easy assembly
- +/- environment sensitive

Find optimal material (combination)

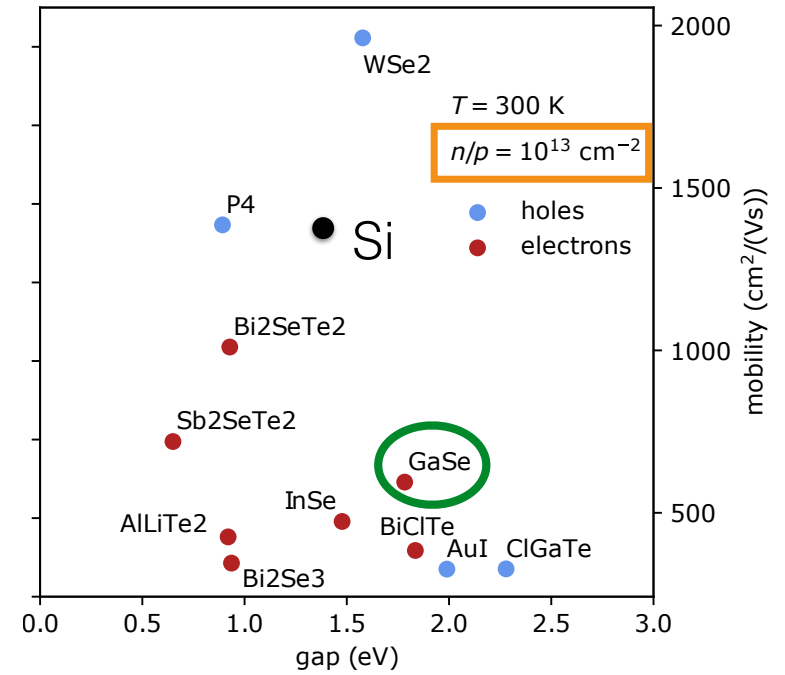
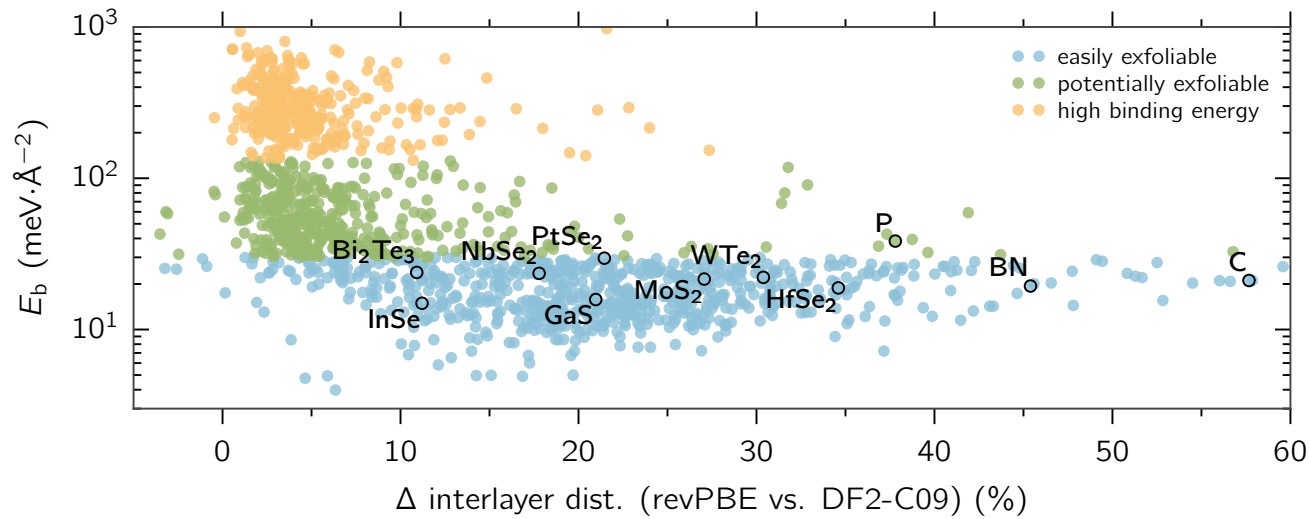


3D Si FinFET
Samsung (2020)



2D graphene
Chen IBM (2008)

Filter the interesting candidates



100 000 3D parents

2000 easily exfoliable

250 small unit cells

12 mobility calculation

Mounet Nat. Nano 13 246 (2018)

Sohier 2D Mater. 8 015025 (2020)



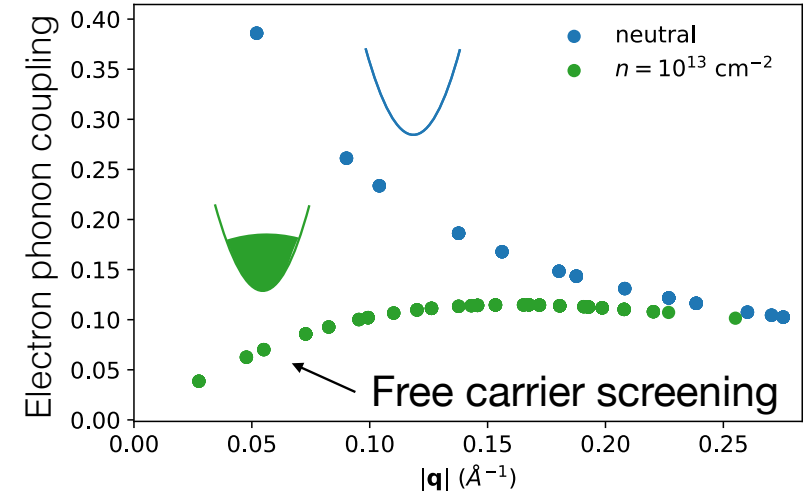
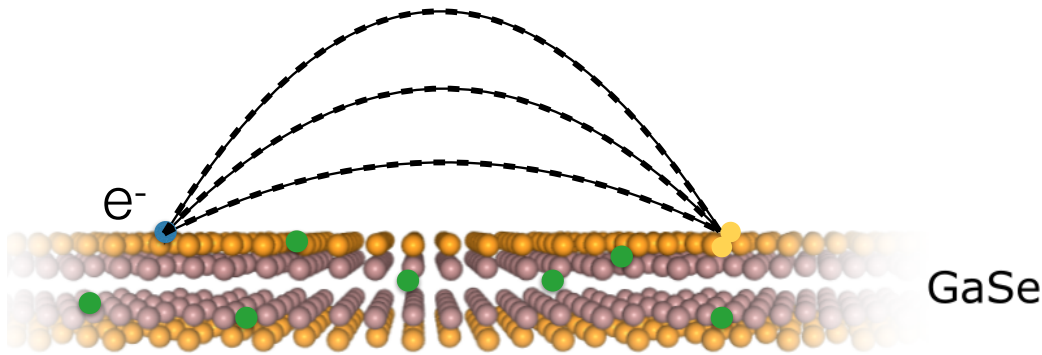
Competitive with Silicon

Good mobility only at high doping

High mobility without doping?

Let's look at GaSe

Transport in 2D



Mobility limited by phonon scattering

Polar mode coupling strongest (Fröhlich)

2D screening geometrical factor $1/|q|$

$$\mathbf{D}(\mathbf{q}) = \frac{e^2}{\Omega} \sum_a \mathbf{Z}_a \cdot \mathbf{u}_{\text{LO}}^a(\mathbf{q})$$

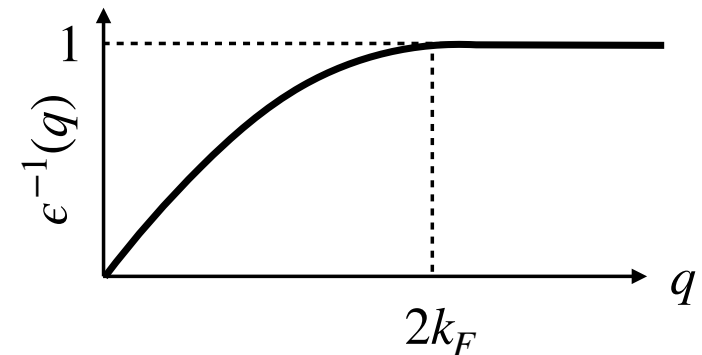
$$\mathbf{E}(\mathbf{q}) = -W_c(\mathbf{q}) [\mathbf{q} \cdot \mathbf{D}(\mathbf{q})] \mathbf{q}$$

$$W_c(\mathbf{q}) = \begin{cases} \frac{4\pi}{|\mathbf{q}|^2 \epsilon} & (3\text{D}) \\ \frac{2\pi}{|\mathbf{q}| \epsilon} & (2\text{D}) \end{cases}$$

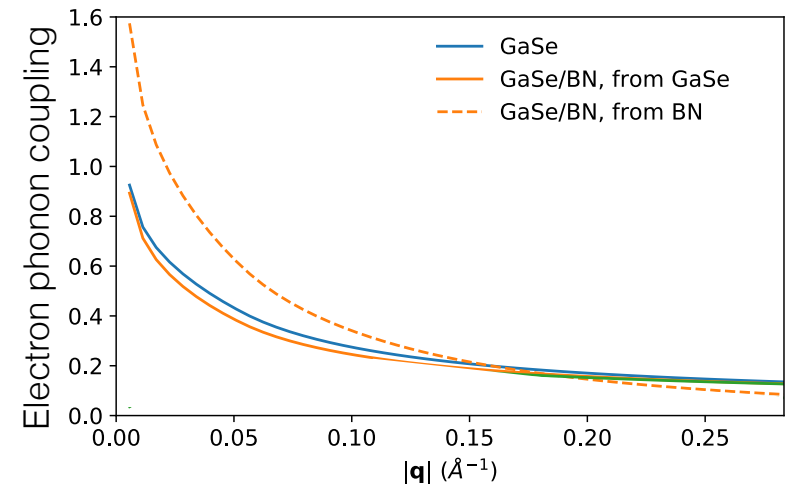
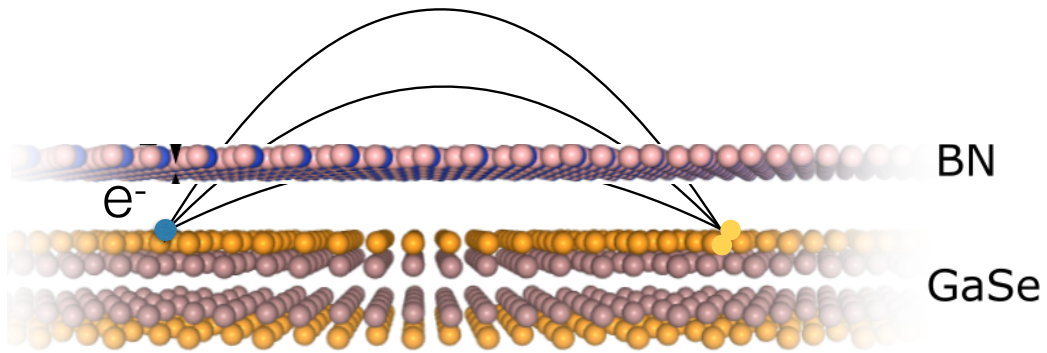
ϵ also depends on phonon wavevector q

Free carriers \rightarrow Drude $\epsilon \rightarrow$ kill EPC

But now it's metallic w/ large voltages



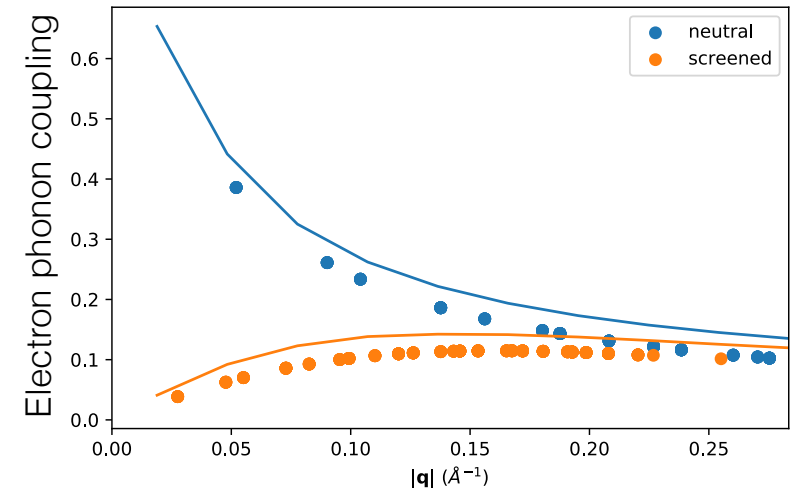
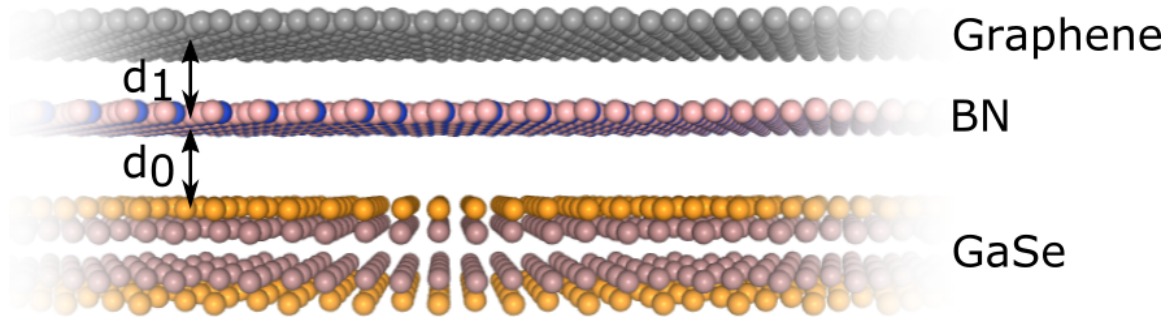
Engineering GaSe



Real device will be encapsulated

Coupling screened remotely with BN (a bit)

Engineering GaSe



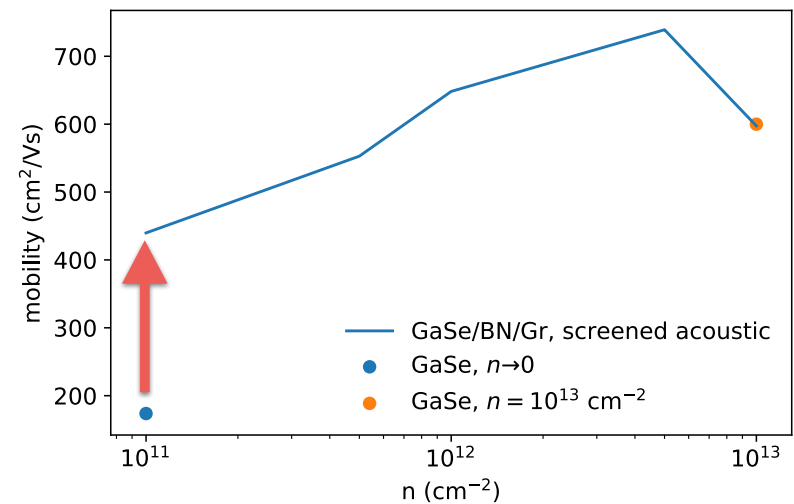
Real device will be encapsulated

Coupling screened remotely with BN (a bit)

Idea: screen remotely with graphene / BN

High mobility also at low doping

→ **New model for heterostructure response**



Sohier Gibertini Verstraete PRMater 5 024004 (2021)

Coupled heterostructure equations

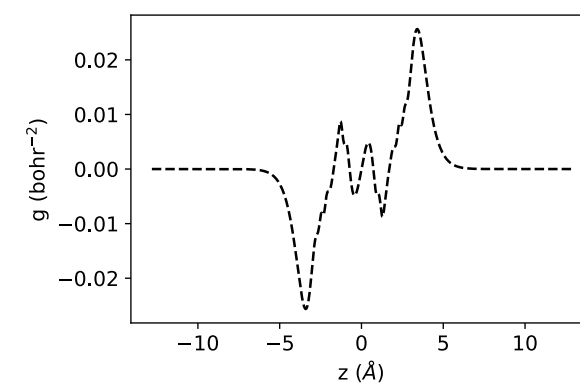
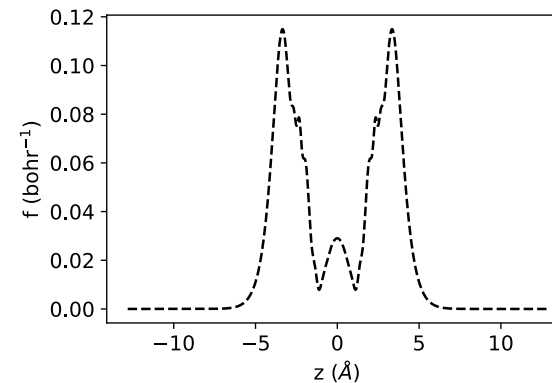
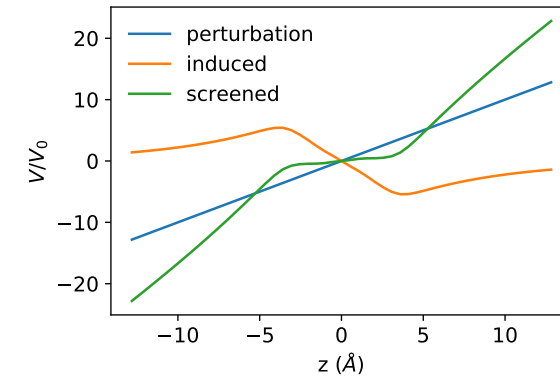
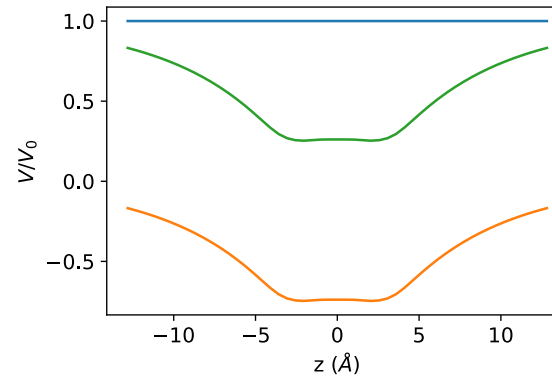
$$\delta n^k(q, z) = \int \chi^k(q, z, z') \left[V_{\text{ext}}(q, z') + \sum_{m \neq k} v_{\text{ind}}^m(q, z') \right] dz'$$

Const / linear perturbation:

$$\chi(q, z, z') = Q(q)f(q, z - z_0)f(q, z' - z_0) + P(q)g(q, z - z_0)g(q, z' - z_0)$$

Full pipeline:

- 1) reference system EPC
- 2) re-screen Fröhlich interaction in layer
- 3) BTE mobility with full band structure

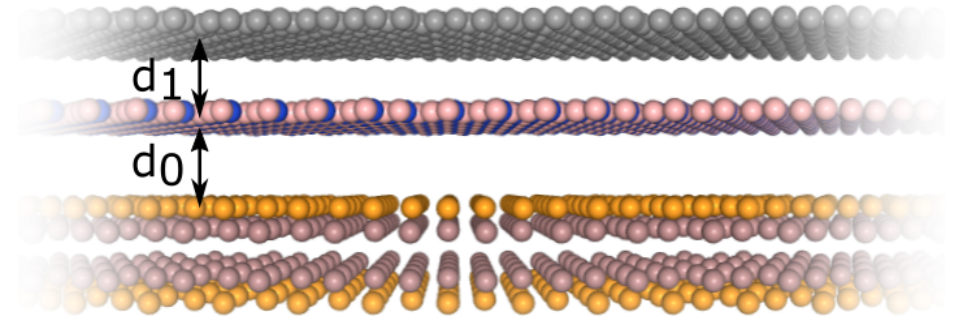


2D transport takeaway

Full stack dielectric model

Any layers any doping

Remote screening concept quantified



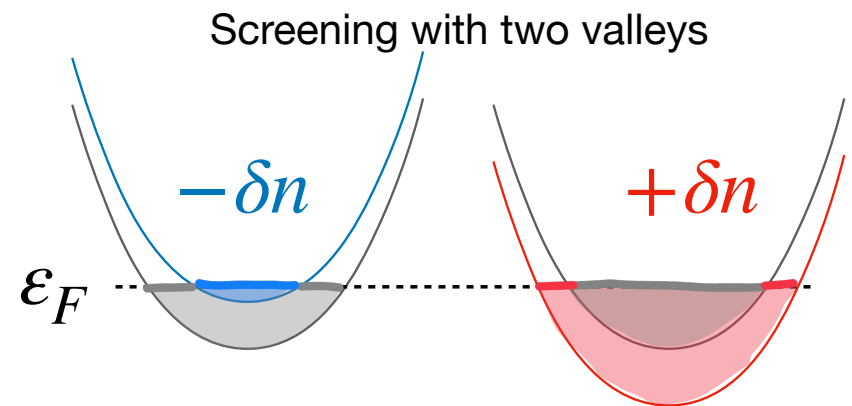
Ongoing work

other phonons: piezo, acoustic

multi valley screening and transport

Sohier Physical Review X (2019)

Sohier Melo Zanolli Verstraete in preparation



Acknowledgements



Universiteit Utrecht



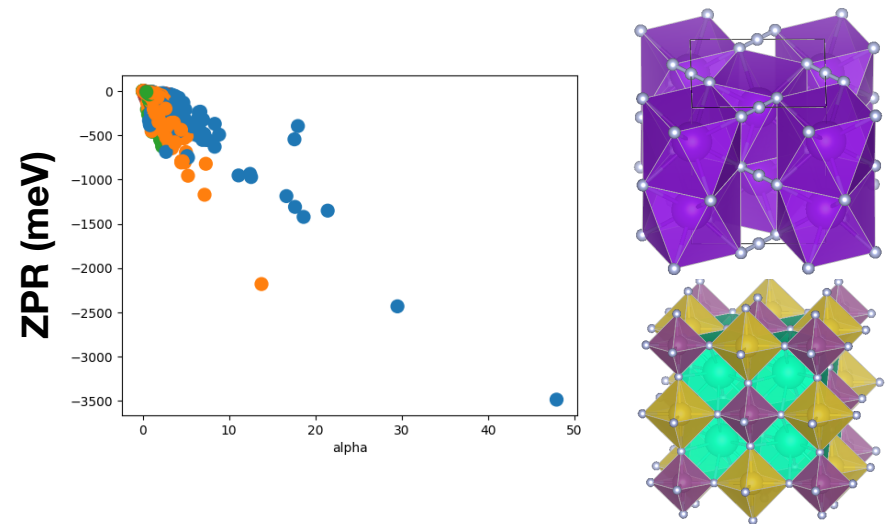
High throughput Fröhlich

Clustering by element period/ionicity

Competing constraints in ϵ_0 ω_{LO} ϵ_∞ m^*

Valid predictor for order of magnitude

Highest LO does not dominate



2D dielectric engineering

Full stack dielectric model

Screen Frohlich w/ remote graphene

Re-calculate EPC and transport

